

Angular Momentum Loss by Magnetic Braking and Gravitational Radiation in Relativistic Binary Stars

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Angular momentum loss (AML) mechanisms and dynamical evolution owing to magnetic braking and gravitational radiation in relativistic binary stars (RBS) are studied with use of physical parameters collected from the literature. We have calculated and compared AML time scales for the RBS with non-degenerate components and double degenerate (DD) systems.

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1. Introduction

White dwarfs (WD), neutron stars (NS), and black holes (BH) represent the endpoint of stellar evolution. Compact objects provide conditions to study physics at high energy, magnetic fields, temperatures etc., and they are dominated by gravitation. Important properties of these compact objects can be determined if they are members of binary systems. This make possible to determine the orbital elements and physical parameters from precise data obtained by various techniques. Such studies are important to test the predictions of stellar evolution theory. Additionally, binary systems with a compact companion or double degenerate (DD), owing to their general relativistic effects, are good laboratories to test general relativity.

The literature has been scanned for the physical parameters of High Mass X-Ray Binary (HMXRB) and Low Mass X-Ray Binary (LMXRB) systems with a black hole (BH) or neutron star (NS) companion, as well as other systems listed in Table 1. We use the data obtained to study the angular momentum loss mechanisms and focus on the dynamical evolution driven by magnetic braking and gravitational radiation in relativistic binaries.

2. Data selection

In modern astrophysics, theoretical and observational studies of compact objects, with both ground based and satellite observations, offer potential to research the formation and evolution. In recent years, as a result of the advanced technology and instrumentation, the number of papers on relativistic binaries has increased. These studies contain significantly accurate orbital elements and physical parameters. The mass interval of these binary systems are summarized in Table 1. Isolated black holes, neutron stars and white dwarfs are included for comparison. Using these parameters collected from the literature we have studied their possible angular momentum loss mechanisms.

3. Angular Momentum Loss

The orbital period of the binary system can be determined accurately. Any orbital period variation in tidally locked binary systems is due to evolutionary changes or the presence of a third body orbiting the binary. In close binaries, in general, we can estimate the angular momentum variation via observations of period changes. Angular momentum loss leads to a period decrease. It can be studied as the orbital angular momentum and spin angular momentum. The spin angular momentum (H^s) for a single star is given in the form,

$$H_i^s = I\omega = k^2 R_i^2 M_i \omega_i, \quad (3.1)$$

where, subscript i stands for the components, $i = 1$ for the primary and $i = 2$ for the secondary, respectively, I is the star's moment of inertia, ω is the spin angular velocity, k^2 is the gyron radius, R_i is the radius of the star, and M_i is the mass of the component. The total orbital angular momentum is

$$H_t^o = M_1 M_2 \sqrt{\frac{Ga(1-e^2)}{M}} = 111q(1+q)^{-2} \left(\frac{M_t}{M_\odot}\right)^{-5/3} \left(\frac{P}{\text{day}}\right)^{1/3} \sqrt{1-e^2} (M_\odot R_\odot^2 / \text{day}) \quad (3.2)$$

Table 1: Mass distribution of relativistic binaries and isolated compact stars based observational data set.

Type	pri+sec	component n	M ₁		M ₂	
			max.	min.	max.	min.
HMXRB	BH+S	6	23.1	6.0	70	6.5
HMXRB	NS+S	11	2.4	1.0	58	8.8
LMXRB	BH+S	13	14	4.0	2.7	0.37
LMXRB	NS+S	6	1.8	1.4	2.3	0.4
NLCV	WD+S/B	39	1.4	0.3	1.1	0.05
mCV	WD+S	10	0.8	0.4	0.5	0.1
PreCV	WD+S	21	0.84	0.39	0.93	0.1
DDNSNS	NS+NS	8	1.6	1.14	1.4	1.05
DDNSWD	NS+WD	9	2.1	1.27	1.3	0.16
DDWDWD	WD+WD	10	0.7	0.32	0.7	0.29
AMCVn	WD+WD	5	0.98	0.59	0.13	0.011
Black Hole		19	23.1	4.0		
Neutron Star		42	2.4	1.0		
White Dwarf		90	1.4	0.01		

where G is gravitational constant, a is the separation of two stars and q is the mass ratio of the components specify $q = \frac{M_2}{M_1}$.

The angular momentum is extracted from the binary orbit because of the (i) magnetic braking in a magnetized stellar wind (MSW), (ii) gravitational radiation (GR), (iii) non-conservative mass transfer and (iv) a third (or multi-) body effect. The total system angular momentum loss is the sum of all these effects. We assume throughout this study that relativistic binary star evolution is driven by GR and MSW.

3.1 Magnetic Braking in a Magnetized Stellar Wind

The Sun is known to rotate more rapidly at the equator than at the poles. The speed of this differential rotational and convection are important for the production of magnetic fields in the convective zone (see Yakut & Eggleton and references therein). Light curves of close binary stars show one maximum higher than the other. This phenomenon is called the O’Connell effect (Milone 1968). These asymmetries are usually attributed to the magnetic activity, though the effect is not satisfactorily explained by existing the models. A large number of late-type stars show magnetic activity indicators. Hence, activity may play an important role in relativistic binaries whose components are late-type stars ($M \lesssim 1.5 M_\odot$).

The Skumanich relation (Skumanich 1972, Smith 1979) predicts a strong correlation between equatorial rotation velocity and age as

$$v_e = 10^{14} \times f \times t_0^{-0.5}$$

cm s⁻¹, where t_0 is the age of the star and the value of f was determined empirically by Skumanich (1972) to be 0.7 and Smith (1979) to be 1.78. The Skumanich law and Eq. (3.1) yield the angular momentum loss by MSW,

$$\left(\frac{dH}{dt}\right)_{\text{MSW}} = -1.6 \times 10^{-30} M_2 R_2^4 \omega^3. \quad (3.3)$$

The angular momentum loss time scale

$$\tau_{MSW} = 14 \left(\frac{M}{M_{\odot}} \right)^{2/3} \left(\frac{R_2}{R_{\odot}} \right)^{-4} \left(\frac{P}{\text{day}} \right)^{10/3} (1+q)^{-1} (1-e^2)^{1/2} \times \text{Gyr}. \quad (3.4)$$

If the secondary stars in binaries fill their Roche lobes then the radius may be calculated with the relation given by Eggleton (1983),

$$R_2 = R_L = a \frac{0.49q^{2/3}}{0.6q^{2/3} + \ln(1+q^{1/3})}. \quad (3.5)$$

Otherwise we used the mass-radius relations

$$R = M^{0.97} \text{ for } M_2 \leq 0.5M_{\odot} \text{ and } R = M^{1.11} \text{ for } 0.5M_{\odot} < M_2 \leq 1.5M_{\odot} \quad (3.6)$$

given by Yakut et al. (2008a).

3.2 Gravitational Radiation

Binary star systems, in which both components are white dwarfs, neutron stars, or black holes emit gravitational waves that can carry off angular momentum. The shape of the binary orbit and the energy of the system given by

$$r = \frac{a(e^2 - 1)}{1 + e \cos \nu} \quad \text{and} \quad E = \frac{-GM_1M_2}{2a}. \quad (3.7)$$

These equations and the semi-major axis variation rates (Landau & Lifshitz, 1951) yield the angular momentum loss by the gravitational radiation and the angular momentum loss time scale (in Gyr)

$$\frac{da}{dt} = -\frac{64}{5} \frac{G^3}{c^5} \frac{M_1M_2}{a^3} (M_1 + M_2), \quad (3.8)$$

$$\left(\frac{dH}{dt} \right)_{GR} = -3.44 \times 10^{44} (M_1M_2M^{-1/3})^2 P^{-7/3} f(e) \quad (3.9)$$

$$\tau_{GR} = 376.4q^{-1} (1+q)^2 M^{-5/3} P^{8/3} (1-e^2)^{3/2} \left(1 + \frac{7}{4}e^2\right)^{-1} \text{Gyr}. \quad (3.10)$$

3.3 Timescales

The angular momentum loss time scales for relativistic and non-relativistic binary stars (LTCB and HTCB) are shown in Figs.1 and 2. The details will be discussed in Yakut et al. (2008b). Fig. 3 is the same as Figs.1 and 2 but for double compact degenerate binary systems.

4. Conclusion

The aim is to collate the physical properties of the relativistic binary stars because they are valuable in studies of formation and evolution of compact objects. From observational data (see Table 1) we also summarize mass distribution for the compact objects. According to the data set, maximum value for mass of black holes, neutron stars and white dwarfs are 23, 2.4 and 1.4

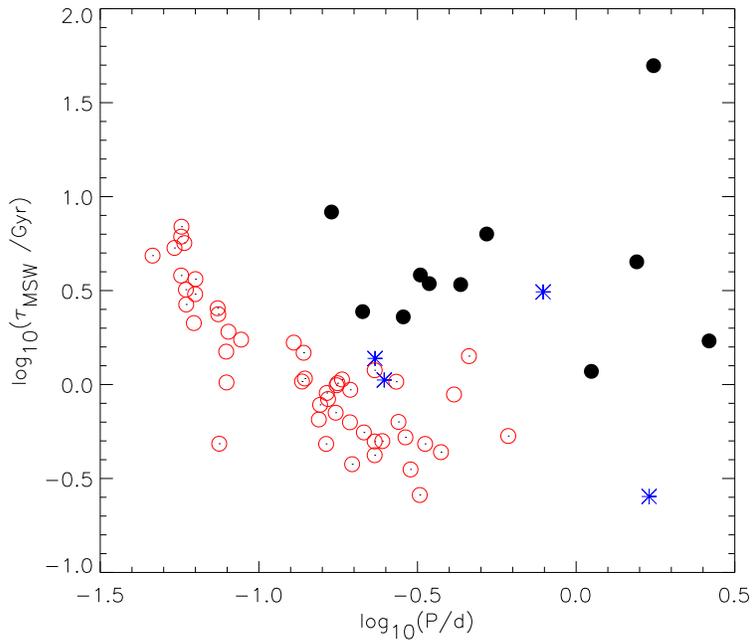


Figure 1: Plot of $\log \tau_{\text{MSW}}$ vs. $\log P$. The open circles, asterisks, and filled circles show binaries with WD, NS and BH companions, respectively.

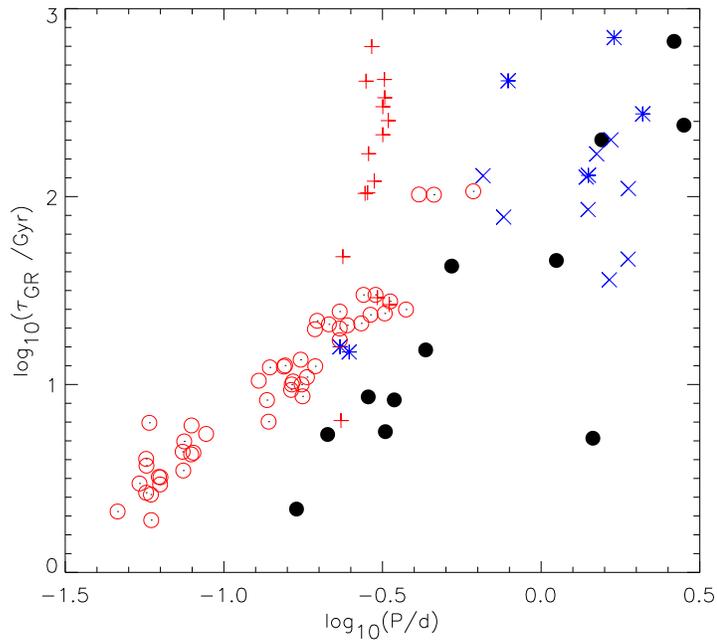


Figure 2: Plot of $\log \tau_{\text{GR}}$ vs. $\log P$. The open circles, asterisks, filled circles show binaries with WD, NS, and BH companions and + and x show the LTCB and HTCB systems, respectively.

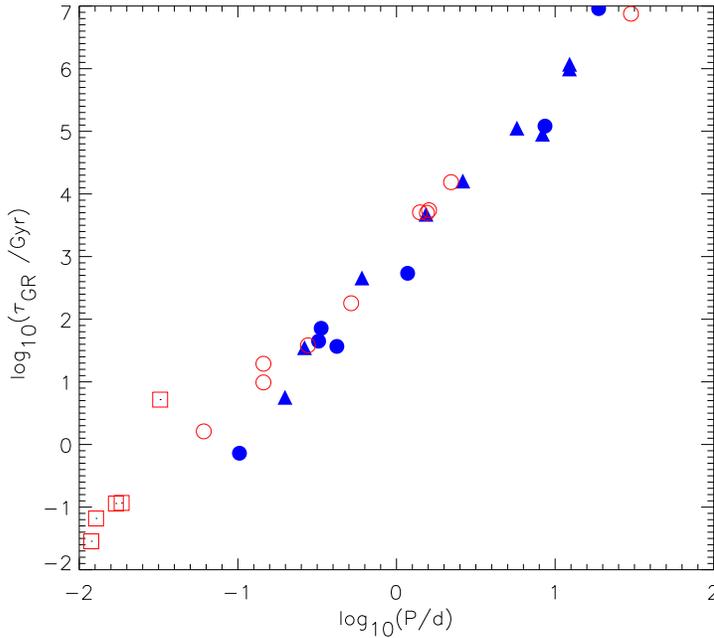


Figure 3: Plot of $\log \tau_{\text{GR}}$ vs. $\log P$ for double degenerate binaries. The open circles, open squares, filled circles and triangles indicate WD+WD, WD+WD (AMCVn), NS+NS and NS+WD binaries, respectively.

M_{\odot} , respectively and for minimum mass for black holes, neutron stars and white dwarfs are 4, 1, and 0.01 respectively. We also draw attention to gap between minimum black hole masses and maximum mass of neutron stars (see for details Yakut et al. 2008b). The data obtained have been used to study the angular momentum loss and its evolution driven by magnetic braking and gravitational radiation. Binary stars with BH components have longer orbital periods than those with WD and NS components. Increase in the orbital period increases the efficiency of AML by GR. Greater the difference between the mass ratio the shorter is the AML time-scale via MSW.

Acknowledgments

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