



Correction to: “Curved space and particle physics effects on the formation of Bose–Einstein condensation around a Reissner–Nordström black hole”

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We have noticed that $\frac{-18}{5}$ in γ_2 in Eq. (24) in our paper [1] should be replaced by $\frac{-9}{5}$. After correction, γ_2 is replaced by

$$\gamma_2 = 2 \left[\frac{-9}{5} + \frac{63}{40} \left(\frac{M}{Q} \right)^2 \right] M. \quad (1)$$

Hence, Eqs. (25)–(27) in [1] are replaced by

$$(m^2 - \omega^2) - \frac{M^2}{9Q^4} \left(\frac{9}{5} - \frac{63}{40} \frac{M^2}{Q^2} \right)^2 = 0, \quad (2)$$

$$- \frac{3M^2}{Q^2} \left(\frac{9}{5} - \frac{63}{40} \frac{M^2}{Q^2} \right) + \left(\frac{9}{5} - \frac{9}{20} \frac{M^2}{Q^2} \right) \left(\frac{4}{5} - \frac{9}{20} \frac{M^2}{Q^2} \right) + (q^2 - m^2)Q^2 = 0, \quad (3)$$

$$\frac{2M}{3Q^2} \left(\frac{9}{5} - \frac{63}{40} \frac{M^2}{Q^2} \right) \left(\frac{9}{5} - \frac{9}{20} \frac{M^2}{Q^2} \right) - 2(qQ\omega - m^2M) = 0. \quad (4)$$

Accordingly, the solutions (29)–(33) for m , ω , and q given in the paper [1] get the following revised forms:

$$m_1^2 = - \frac{9(7M^4 + 2M^2Q^2)}{40Q^6}, \quad m_2^2 = \frac{9(7M^2 - 8Q^2)^2(4Q^2 - 3M^2)}{1600Q^6(Q^2 - M^2)}, \quad (5)$$

$$\omega_{1,2} = \mp \frac{3iM(7M^2 + 12Q^2)}{40Q^4}, \quad \omega_{3,4} = \mp \frac{3i(7M^4 - 22M^2Q^2 + 16Q^4)}{40\sqrt{Q^8(M^2 - Q^2)}}, \quad (6)$$

$$q_{1,2} = \mp \frac{3i(17M^2 - 8Q^2)}{20Q^3}, \quad q_{3,4} = \pm \frac{3iMQ(27M^2 - 28Q^2)}{40\sqrt{Q^8(M^2 - Q^2)}}, \quad (7)$$

where $\{m_1^2, \omega_{1,2}, q_{1,2}\}$ and $\{m_2^2, \omega_{3,4}, q_{3,4}\}$ are two different sets of solutions. We notice that $q_{1,2}$ are imaginary, so $\{m_1^2, \omega_{1,2}, q_{1,2}\}$ are not physical solutions. Therefore, only the set $\{m_2^2, \omega_{3,4}, q_{3,4}\}$ corresponds to physical solutions with the requirement $|Q| > M$. It is evident that these solutions cover a much wider ranges of m^2 , ω , q than those in [1]. The new ranges are $1.44 < m^2Q^2 < \infty$, $1.2 < \omega Q < \infty$, $0 < q Q < \infty$ where m^2Q^2 and ωQ are in the order of 1 for most of the values of M and Q . $q Q$ is in the order 10^{-1} for most of the values of M and Q .

Accordingly, ψ_ω in (28) in [1] gets the following revised form:

$$\psi_\omega = \exp \left\{ \left(\frac{\pm 3M(7M^2 - 8Q^2)}{40Q^4\sqrt{Q^2 - M^2}} \right) \left[(2M^2 - Q^2) \left(\tan^{-1} \left(\frac{M - r}{\sqrt{Q^2 - M^2}} \right) - \tan^{-1} \left(\frac{M - r_0}{\sqrt{Q^2 - M^2}} \right) \right) \right. \right. \\ \left. \left. - \sqrt{Q^2 - M^2} \left(M \log(-2Mr + Q^2 + r^2) - M \log(-2Mr_0 + Q^2 + r_0^2) + r - r_0 \right) \right] \right\}$$

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Fig. 1 $|\psi_\omega|$ versus \bar{r} graph for $\bar{M} = 10, \bar{Q} = 17, \bar{r}_0 = 20$

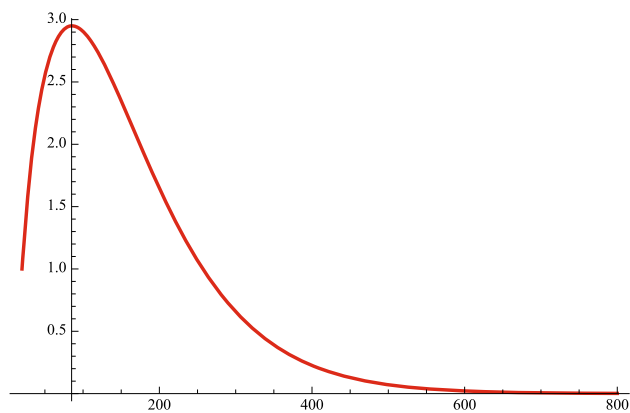
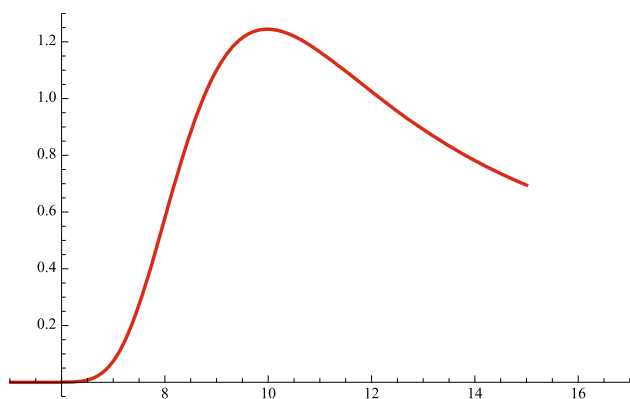


Fig. 2 $|\psi_\omega|$ versus \bar{Q} graph for $\bar{M} = 10, \bar{r} = 20, \bar{r}_0 = 30$



$$\begin{aligned}
 & + \left(\frac{3(6M^3 - 4MQ^2)}{40Q^2\sqrt{Q^2 - M^2}} \right) \left(\tan^{-1} \left(\frac{M - r}{\sqrt{Q^2 - M^2}} \right) - \tan^{-1} \left(\frac{M - r_0}{\sqrt{Q^2 - M^2}} \right) \right) \\
 & - \frac{3(3M^2 + 8Q^2)}{40Q^2} \left(\log(-2Mr + Q^2 + r^2) - \log(-2Mr_0 + Q^2 + r_0^2) \right) \\
 & + 3 \log(r) - 3 \log(r_0) \Big\}, \tag{8}
 \end{aligned}$$

After this correction Fig. 1 in [1] is replaced by Fig. 1 above which is essentially the same as the one in [1]. Figure 2 in [1] now becomes irrelevant. Figure 3 in [1] is replaced by Fig. 2 above which is essentially the same as the one in [1].

The essential difference between the original paper and the one after the correction is: The case $|Q| < M$ is excluded after the correction. The allowed ranges of the parameters m^2, ω, q after correction for $|Q| > M$ are much wider than those obtained in [1] (and they include the ranges of parameters found in [1] as subcases).

Reference

1. R. Erdem, B. Demirkaya, K. Gültekin, Eur. Phys. J. Plus **136**(9), 972 (2021)